# A Statistical Mechanics Approach to Holographic RG

## Bethe lattice Ising model & p-adic AdS/CFT Kouichi Okunishi / D02 Group Tadashi Takayanagi / CO1 Group









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The Bethe lattice Ising model can be served as a prototype toy model for the holographic RG ⇒ p-adic AdS/CFT



#### Mean field approximation





all to all interactions (Husimi-Temperley model) Bethe approximation



# Tensor network and holography

The MERA(1+1D) describes the area law of entanglement entropy up to log correction.

Ryu-Takayanagi formula (AdS side)

### However

- tensor network states are usually variational wavefunction (CFT side)
- MERA tensor elements are highly nontrivial because of numerical optimization

Is there any simple tensor network enabling us to analytically describe the AdS side?

# Tensor network---MERA/TNR



MERA network representation



Variational wavefunction (CFT side)

The EE between a system and the rest

 $\gamma_A \sim \#$  of bonds  $S_{\rm EE} \sim \log \ell$ 





## types of TNs



## Cayley tree Ising model

A tree tensor network (loop free) analytic calculation! (1D model) The center spin = Bethe Approximation (Mean field)



$$Z_N = \sum_{\sigma_0} e^{h\sigma_0} [g_N(\sigma_0)]^q$$
$$g_N(\sigma_0) = \sum_{\sigma_1} e^{K\sigma_0\sigma_1} [g_{N-1}(\sigma_1)]^p$$
$$= \sum_{\sigma_1} e^{K\sigma_0\sigma_1} \left[ \sum_{\sigma_2} e^{K\sigma_1\sigma_2} [g_{N-2}(\sigma_2)]^p \right]$$

T.P. Eggarter Phys. Rev. B 9 2989 (1974)E. Mullar-Hartmann and J. Zitterz , PRL 33 893 (1974)R. J. Baxter, textbook, chap. 4



Recursion relation for the effective magnetic field

$$C_{n-1}e^{\tilde{h}_{n-1}\sigma_{n-1}} = \left[\sum_{\sigma_n} e^{K\sigma_{n-1}\sigma_n + \tilde{h}_n\sigma_n}\right]^p$$
  

$$\tilde{h}_{n-1} = f(\tilde{h}_n)$$
  
with  $f(x) \equiv \frac{p}{2}\log\left(\frac{e^{K+x} + e^{-K-x}}{e^{K-x} + e^{-K+x}}\right)$ 

 $ilde{h}_N\equiv h_b~~$  boundary magnetic field (no magnetic field in the bulk )

This can be viewed as a holographic RG flow from the boundary to the Bulk

Bethe Approximation 
$$ilde{h}^* = f( ilde{h}^*)$$

Self-consistent equation = fixed point of RG flows

### Holographic RG flow from the boundary to the bulk



Effective magnetic field around the fixed point

$$\tilde{h}_n \simeq h_b \lambda^{N-n} \sim h_b \lambda^{-n}$$
 with  $\lambda = p \tanh K$ 

**Correlation function** for boundary spins ( $h_b = 0$ )

$$\langle \sigma_{Nk} \sigma_{Nk'} \rangle = (\tanh K)^{2d_{kk'}}$$

 $d_{kk^\prime} = N - n_{kk^\prime}$  specifies the lattice distance for two spins

For the redline path N = 4, k = 0, k' = 7 $n_{kk'} = 1 \text{ and } d_{kk'} = 3$ 

The both exhibits the exponential decay at the level of the lattice network!

## The boundary spin operator for $h_b = 0$

. . .

$$\sum_{\sigma_0 \cdots \sigma_N} \sigma_N e^{K \sum_{i=0}^{N-1} \sigma_i \sigma_{i+1}}$$
$$= 2 \cosh K \sum_{\sigma_0 \cdots \sigma_{N-1}} [\tanh K \sigma_{N-1}] e^{K \sum_{i=0}^{N-2} \sigma_i \sigma_{i+1}}$$

$$= (2\cosh K)^{N-n} \sum_{\sigma_0 \cdots \sigma_n} [(\tanh K)^{N-n} \sigma_n] e^{K \sum_{i=0}^{N-n+1} \sigma_i \sigma_{i+1}}$$

$$(\tanh K)^{N-n}\sigma_n = \sigma_N$$

consistent with the boundary correlation function

## network geometry



**Radial direction** 

$$R_n = 1 - z_n$$
$$z_n = p^{-n}$$

**Circumference direction** 

$$x_k = \frac{2\pi R_n}{qp^{n-1}}k \simeq \frac{2\pi p}{q}z_n k,$$

for 
$$n \gg 1$$

$$\tilde{h}_n \sim \lambda^{-n} = z_n^{1-\Delta}$$

$$\langle \sigma_{N0}\sigma_{Nk'}\rangle \sim \frac{1}{(x_{k'}-x_0)^{2\Delta}}$$

$$\Delta \equiv - \tfrac{\log(\tanh K)}{\log p}$$



Unit disk representation  

$$(z_n, x_k)$$
 $ds^2 = \frac{L^2}{z^2} (dz^2 + d\tilde{x}^2)$   
near the boundary
 $L \equiv \frac{1}{\log p}$ 
A scalar field in AdS  
 $\phi(z) \sim Az^{\Delta_-} + Bz^{\Delta_+}$  with  $\Delta_+ + \Delta_- = d$   
 $A \leftrightarrow h_b$   
 $B \leftrightarrow \sigma_N$  with  $d = 1$   
 $\Delta_+ \leftrightarrow \Delta$   
 $\Delta \equiv -\frac{\log(\tanh K)}{\log p}$ 

Notes:

\*The correlation path can be described by a geodesic in Poincare coordinate

\* EE for the Bethe lattice is bounded because of the tree network nature.



Be the transition point corresponds to  $\Delta=1$  (massless point)

$$\Delta \equiv -\frac{\log(\tanh K)}{\log p} \qquad \qquad \Delta' = -\frac{\log\left(\frac{\sinh(2K)}{\cosh(2K) + \cosh(2\tilde{h}_*)}\right)}{\log p}$$

T< T\_B



### Ferromagnetic fixed point and crossover for $T < T_B$



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